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Tapered Finite Elements in the Optimality Criterion Approach

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Introduction

THE optimality criterion approach for obtaining minimum mass (volume if density is constant) design of a structure subjected to a variety of constraints is well established.¹⁻⁵ In all earlier studies element detail in the form of area or thickness (constant for an element) is used as a design variable and hence the optimum structure obtained is a stepped one with discontinuities at the element boundaries.

In this Note an alternate method, in which the nodal detail is in the form of nodal area or thickness, is used as a design variable. As a consequence of this the finite elements used to model the structure will have a tapered configuration. This will enable continuity of structure at the element boundaries.

This method is applied to the problem of optimization of a cooling fin with a temperature constraint. This problem has already been solved by a continuum approach⁶ and the optimality criterion approach using constant area elements.⁴ The solution obtained with nodal areas as design variables, when compared to the exact solution,⁶ is more accurate when compared to that obtained with element areas as design variables.⁴

Formulation

Let the cooling fin be divided into M number of linearly tapered finite elements of equal length. Then the volume of the bar V is given by

$$V = \sum_{i=1}^n \frac{A_i l}{p} \quad (1)$$

where n is the number of nodal points, l is the length of the element, and A_i , the design variable, is the area of the fin at node i . The value of p is 2 for $i=1$ or n and 1 for $i=2$ to $(n-1)$.

The constraint equation is

$$T_m - T_c = 0 \quad (2)$$

or

$$\{T\}^T \{Q\} - T_c = 0 \quad (3)$$

where T_m is the temperature at the node m at which the value is constrained to be T_c , $\{T\}$ is the temperature vector of the fin obtained after performing thermal analysis, and $\{Q\}$ is a vector of zeros except at the constraint degree-of-freedom where the value is unity.

Following the procedure explained in Ref. 4, the optimality criterion for the system can be obtained as

$$e_i p / l = 1 \quad (4)$$

where

$$e_i = \lambda \{T\}^T \frac{\partial [D]}{\partial A_i} \{s\} \quad (5)$$

λ is the Lagrangian multiplier, $[D]$ is the sum of the assembled conduction and convection matrices, and $\{s\}$ is the temperature vector obtained with $\{Q\}$ as the thermal load vector.

Recurrence relations for design variables A_i and λ are obtained as⁴

$$A_{j+1} = \sqrt{(e_i p / l)} A_{j_i} \quad (6)$$

and

$$\lambda_{j+1} = \sqrt{T_m / T_c} \lambda_j \quad (7)$$

where j is the iteration number.

It is to be noted here that $[D]$ is of the form

$$[D] = [C] + [H] \quad (8)$$

where $[C]$ is the assembled conduction matrix dependent on A_i and $[H]$ is the assembled convection matrix independent of A_i . Therefore,

$$\frac{\partial [D]}{\partial A_i} = \frac{\partial [C]}{\partial A_i} \quad (9)$$

Also, for a linearly tapered k th finite element, the element matrix $[C]_k$ can be written as,

$$[C]_k = A_k [C_1] + A_{k+1} [C_2] \quad (10)$$

and from Eq. (10), the derivative $\partial [C] / \partial A_i$ can be obtained to be used in Eq. (6).

Numerical Results

The method proposed in this note is applied to the problem of obtaining the minimum mass (volume if density is constant) design of a cooling fin (Fig. 1) subjected to a temperature constraint.

Heat is generated in the fin at a constant rate of q per unit length with convective heat loss over both the top and bottom surfaces. The end $\bar{x}=0$ ($\bar{x}=x/L$, where x is the axial coordinate and L is the length of the fin) is maintained at a temperature of $\bar{T}=0$ ($\bar{T}=T/T_o$, where T_o is the reference temperature) and the insulated end is constrained to a value of $\bar{T}_c=0.3$.

One-dimensional linearly tapered finite elements with cubic temperature distribution over the element are used to model the cooling fin. The elements have two nodes with \bar{T} and $d\bar{T}/d\bar{x}$ as nodal degrees-of-freedom. In the present analysis the fin is idealized with five linearly tapered elements.

To initiate the iterative process an initial design of $\bar{A}_i = 1$, $i=1$ to n (a fin of constant area), and $\lambda = 1$ is used. The

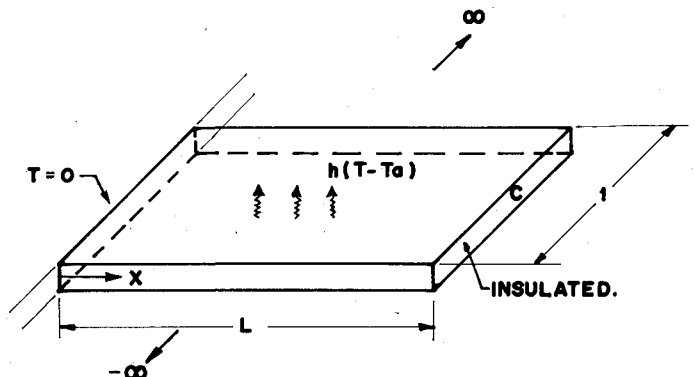


Fig. 1 Fin geometry and boundary conditions.

Table 1 Optimum area (\bar{A}) and temperature (\bar{T}) distribution along length of the cooling fin – five-element solution ($\bar{T}_c = 0.3, N_G = 1.0, B_i = 0.0$)

Node	\bar{A}			\bar{T}		
	Present method	Exact solution ⁶	Ref. 4	Present method	Exact solution ⁶	Ref. 4
1 ^a	2.2263 ^c	2.2222	1.9464 ^d	0.0	0.0	0.0
2	1.9915	1.9876	1.7846	0.0853	0.0853	0.0930
3	1.7281	1.7213	1.5938	0.1604	0.1650	0.1715
4	1.4092	1.4054	1.3377	0.2239	0.2241	0.2343
5	1.0271	0.9938	0.9213	0.2727	0.2732	0.2790
6 ^b	0.3055	0.0		0.3000	0.3000	0.3000
\bar{V}_{opt}	1.4845	1.4815	1.5168			

^aNode at $\bar{T} = 0$. ^bNode at insulated end. ^cNodal areas of linearly tapered elements (solution at 25 iterations). ^dConstant area elements (solution at 20 iterations).

Table 2 Effect of convection on minimum volume for cooling fin – five-element solution ($\bar{T}_c = 0.3, \bar{T}_a = 0.05, N_G = 1.0$)

B_i	\bar{V}	
	Present method	Constant area elements (Ref. 4) ^c
0.0 ^a	1.4845 (25) ^b	1.5168 (1.5023) ^d
0.05	1.4564 (25)	1.4878 (1.4736)
0.2	1.3712 (30)	1.4005 (1.3870)
0.5	1.1985 (34)	1.2232 (1.2115)

^aFor $B_i = 0.0$, $\bar{V}_{exact} = 1.4815$ (Ref. 6). ^bNumber of iterations to obtain convergence. ^cResults at 20 iterations. ^dNumbers in the brackets represent 20-element solution.

iterative process is terminated when the constraint temperature converges up to four significant figures.

Table 1 gives the nondimensional area distribution \bar{A} ($\bar{A} = A/A_0$, where A_0 is the reference area) and temperature distribution \bar{T} at the six nodes of the fin for the heat generation number, $N_G = 1.0$ ($N_G = qL^2/K\bar{A}_0T_0$, where K is the thermal conductivity of the material), without convective heat loss, i.e., the Biot number, $B_i = 0$ ($B_i = hL^2/K\bar{A}_0$, where h is the convective heat-transfer coefficient). The exact solution for this case⁶ and the solution obtained with constant area elements with element areas as design variables⁴ are also presented in this table. Nondimensional optimum fin volume \bar{V}_{opt} ($\bar{V}_{opt} = V/A_0L$, where V is the volume of the fin) is given in this table. The present solution converged in 25 iterations and the solution of Ref. 4 is given at 20 iterations. It can be seen from Table 1 that the results obtained with nodal areas as design variables agree well with the exact solution, whereas those obtained with element areas as design variables are inferior and give higher fin volume compared to the present solution.

In Table 2, the effect of convective heat loss on the minimum mass design is shown. In this table, the nondimensional volume \bar{V} is presented for various values of B_i (the ambient temperature \bar{T}_a is taken as 0.05). \bar{V} obtained from the present five-element solution is less compared to the 20-element solution of Ref. 4. Also it can be noted here that as B_i increases, the number of iterations required for convergence increases. This is because the convection matrix is independent of design variables. From this table it can be concluded that the effect of convective heat loss is to decrease the optimum volume.

Conclusion

The concept of nodal variable as a design variable in the optimality criterion approach for obtaining minimum mass design of structures is proposed. The effectiveness of this

proposal is shown through an example of obtaining the minimum mass (volume if density is constant) design of a cooling fin subjected to a temperature constraint.

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Amplification of Finite-Amplitude Waves in a Radiating Gas

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Introduction

USING the theory of singular surfaces, a number of problems relating to wave propagation in diverse branches of continuum mechanics have been studied previously.¹⁻⁷ The present paper uses the singular surface theory of Thomas⁶ to study the propagation of arbitrarily shaped finite-amplitude waves in a radiating gas near the optically thin limit. The gas is assumed to be perfect, optically gray, and in thermodynamic equilibrium. It is found that a

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